On using the BRDF for simulations with turbulence

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Active imaging system in turbulence

- R = Bi-static separation
- L = Range
- \hat{n} = Normal vector of the target surface



- The laser emits rays towards the target
- These rays are deviated by turbulence.
- They are reflected or scattered off the target surface.
- They then propagate through turbulence towards the imager.



Bi-directional Reflectance Distribution Function

- The BRDF is the distribution of reflected/scattered rays (on the right) for a given incident ray (on the left).
- The red ray corresponds to reflection.
- Whereas the blue rays represent scattering.





The paraxial propagation approximation

- We assume that propagation occurs principally along the z-axis
- The propagator for the electric field is $U = e^{ikz}u$, where k is the wavenumber and u obeys the paraxial equation: $2ik\partial_z u + \nabla^2 u + 2k^2n_1u = 0$, where $n_1 \ll 1$ is the refractive index fluctuation.
- We assume that the Rytov approximation holds so that $u_0 = u \exp[\chi + iS]$.
- We will use the MCF of the field: $\Gamma(\vec{R}, \vec{\Delta}, z) = u(\vec{R} + \frac{1}{2}\vec{\Delta}, z)u^*(\vec{R} \frac{1}{2}\vec{\Delta}, z)$.
- The Wigner function is: $W(\vec{R}, \vec{P}, z) = \int d^2 \Delta \Gamma(\vec{R}, \vec{\Delta}, z) \exp[-ik\vec{P} \cdot \vec{\Delta}]$, where \vec{P} is the slope of the ray.



The BRDF equation

- We assume that the outgoing Wigner function is related to the incident Wigner function by an integral transform.
- $W_o(\vec{R}, \vec{P}, 0) = \int d^2 P' W_i(\vec{R}, \vec{P}', 0) T(\vec{P}, \vec{P}')$
- For the BRDF, we assume a Gaussian scattering function:

$$T(\vec{P}, \vec{P}') = \frac{1}{2\pi\sqrt{|\Sigma|}} \exp\left[-\frac{1}{2}(P_i + P_i' - 2Q_i)\Sigma_{ij}^{-1}(P_j + P_j' - 2Q_j)\right]$$

• Where the symmetric matrix Σ_{ij} represents the spreads of the scattering function and \vec{Q} is the slope of the normal of the target surface.

Note that the spreads can be a function of the position.



The MCF transfer function

While the previous equations may look daunting, in terms of the MCFs we get the straightforward relation:

•
$$\Gamma_o(\vec{R}, \vec{\Delta}, 0) = \Gamma_i(\vec{R}, -\vec{\Delta}, 0) \exp\left[-\frac{k^2}{2}\Delta_i \Sigma_{ij}\Delta_j + 2ik\vec{Q}\cdot\vec{\Delta}\right]$$

- Note that for a Lambertian (incoherent) surface, we have $\Sigma_{ij} = \sigma^2 \delta_{ij}$ where $\sigma^2 \gg 1$ because the surface scatters over a wide range of angles.
- In that case, the exponential approximates a Dirac delta function $\delta(\overline{\Delta})$ and the surface slope no longer matters.



The active imaging model

- We have adapted the DRDC passive imaging model to an active model and we use the BRDF formalism previously developed.
- We tested it on a uniform grey surface with an isotropic BRDF $\Sigma_{ij} = \sigma^2 \delta_{ij}$ characterized by a length scale $l = 1/k\sigma$ that describes its specularity.
- Range = 2.3 km, Outer scale = 10 m, Inner scale = 6 mm, Cn2 = 5e-14, Wavelength = 4.2 μm, IFOV = 3.83 μrad, Aperture = 18.36 cm, Source = 1 cm.





No turbulence case

Range = 2 km, Outer scale = 2 m, Inner scale = 1 cm, Wavelength = 1.55 μm, IFOV = 6.25 μrad, Aperture = 24 cm, Source = 1 cm.























Evaluation of results

- We have developed a rough model for weakly specular surfaces for active imaging through weak turbulence.
- It requires the coherence lengths of the reflecting surface and its orientation with respect to the line-of-sight.
- It assumes that the slopes of the incoming and outgoing rays are small, along with the slope of the surface normal.
- More work is needed to handles cases where
 - Turbulence is strong,
 - The target is highly specular,
 - The surface normal has a large slope,
 - The illuminating beam creates speckle.

